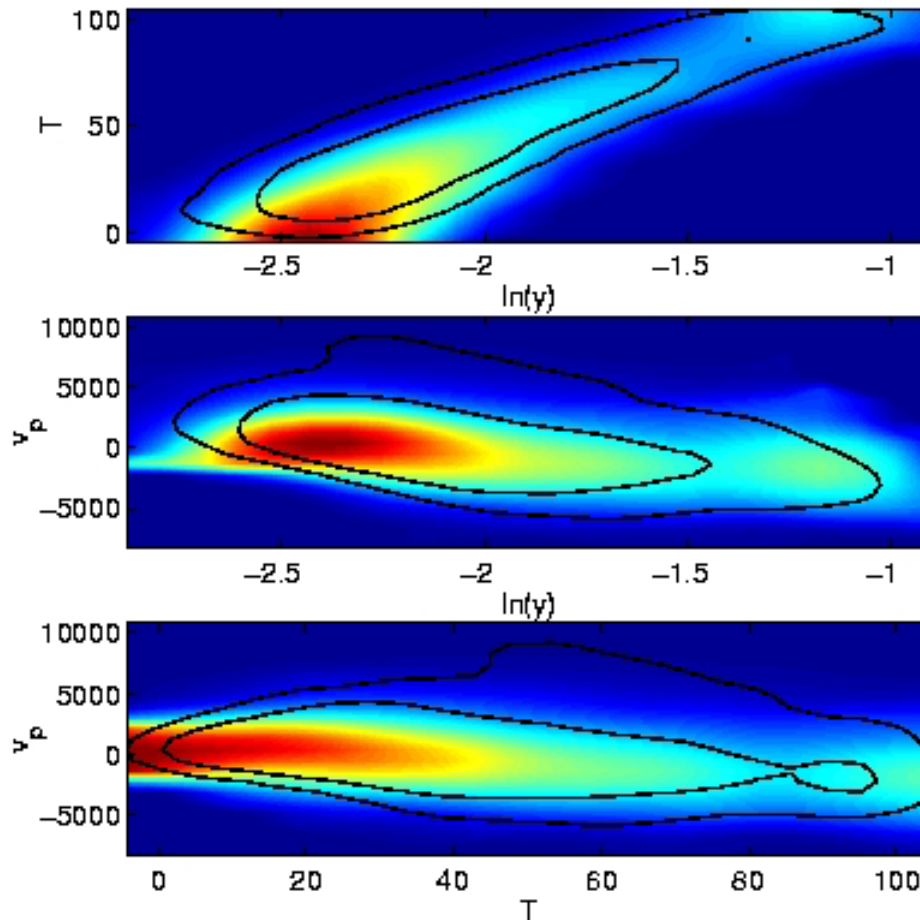


# Cluster Temperatures and SZ observations

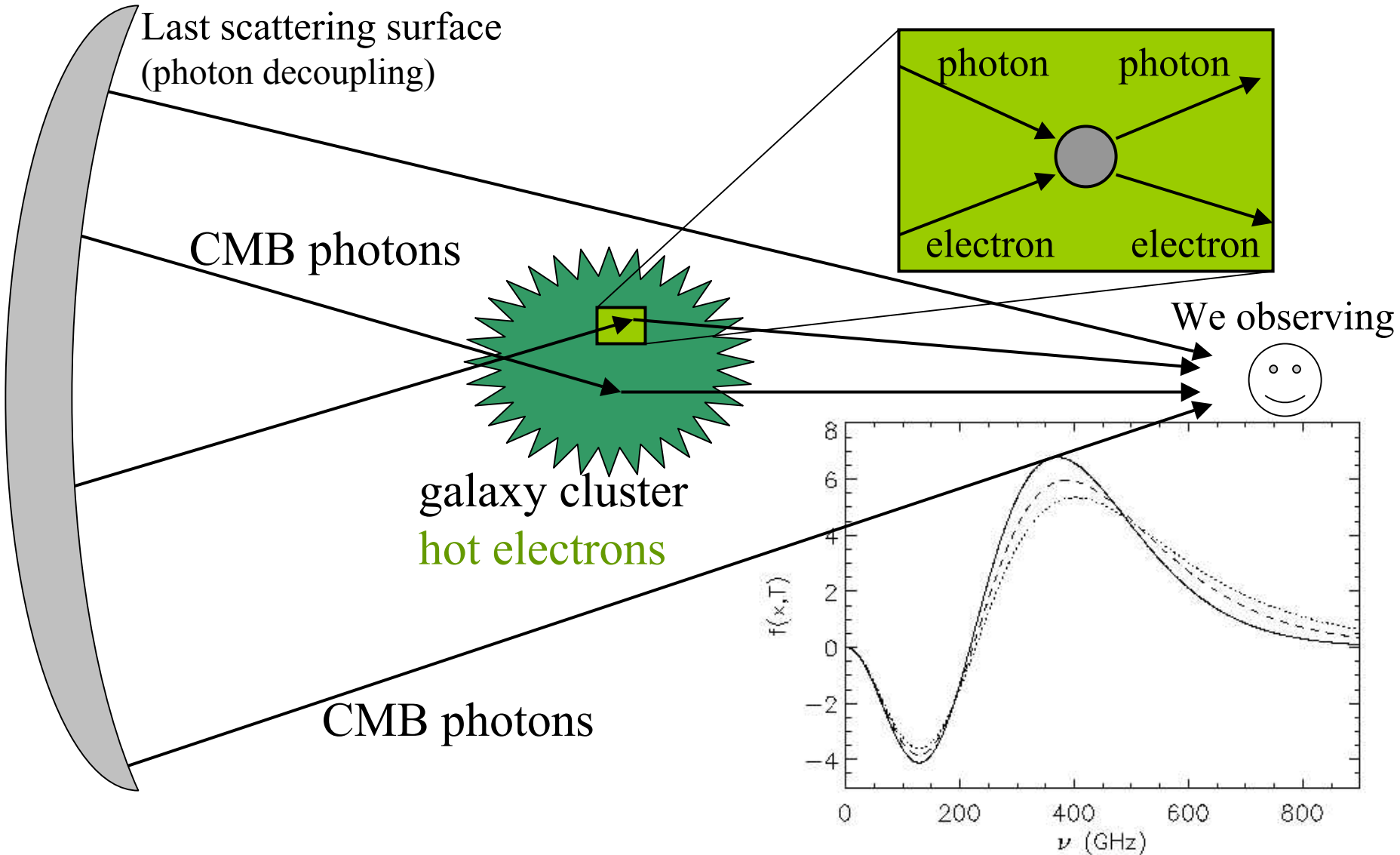
## why we need good angular resolution

IAS, April 2005

Steen H. Hansen, Zurich

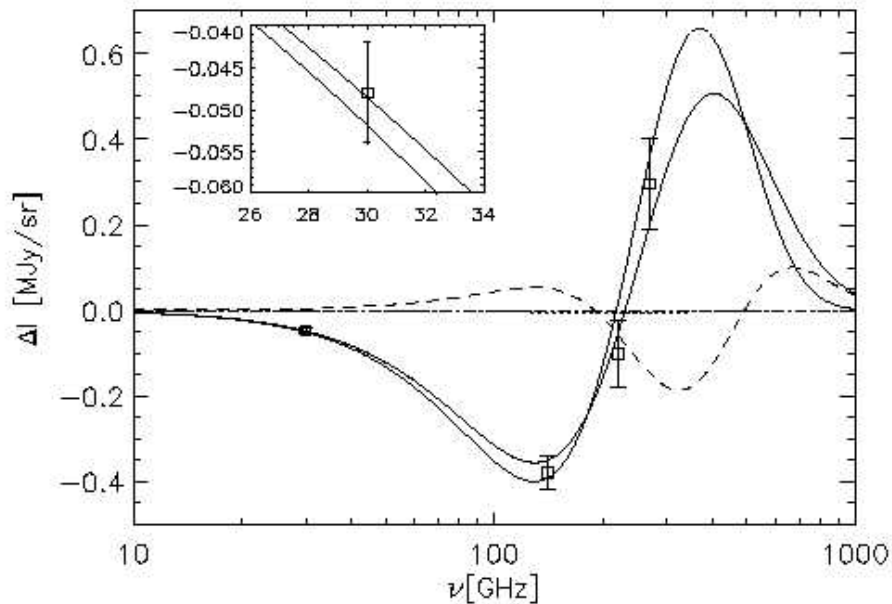


# The SZ effect depends on temperature

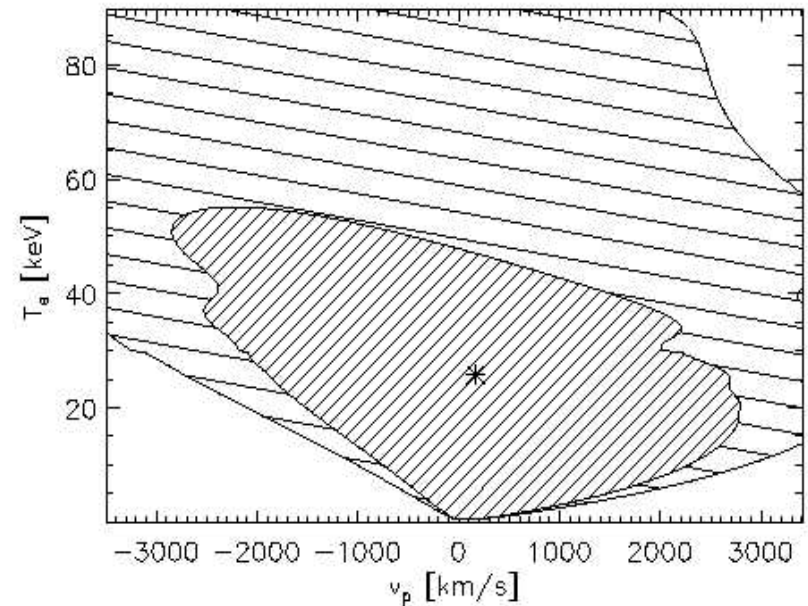


# Measure cluster temperature with SZ

A2163, result is  $T < 60$  keV at  $1\sigma$

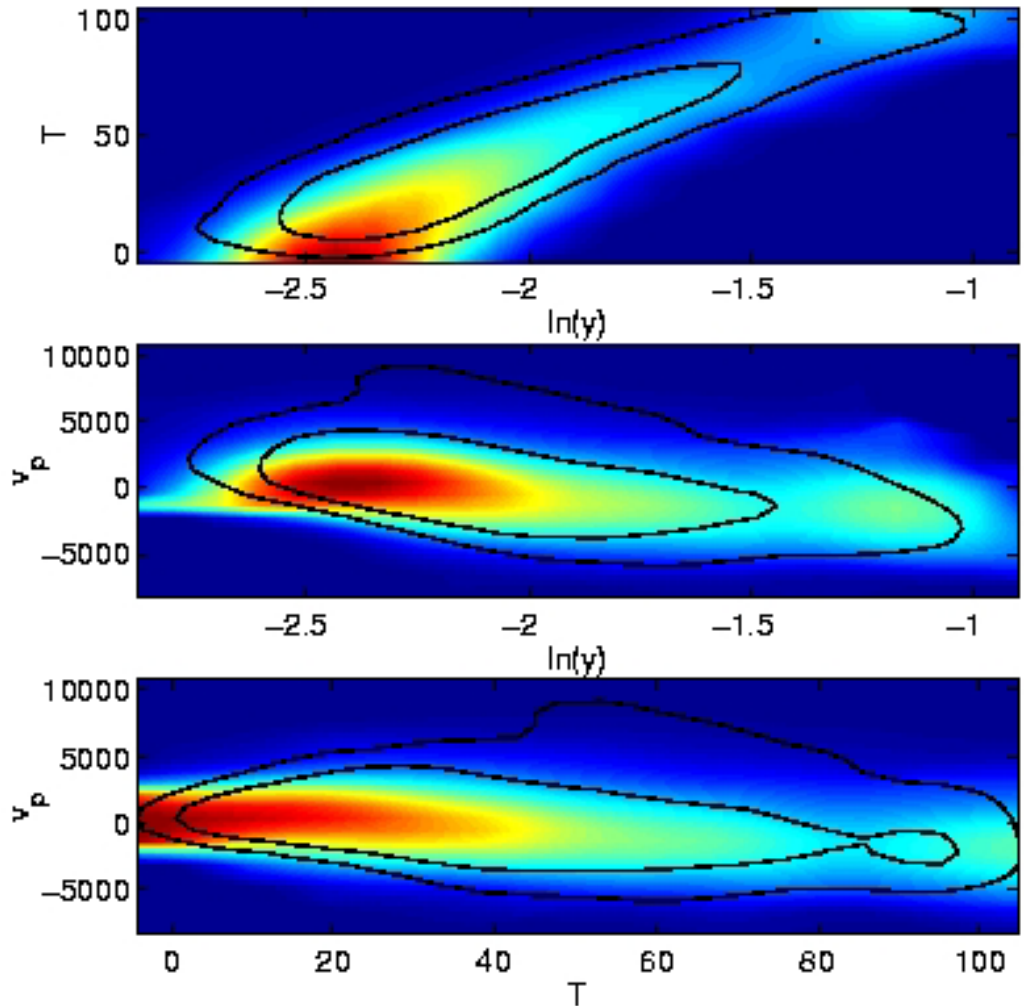
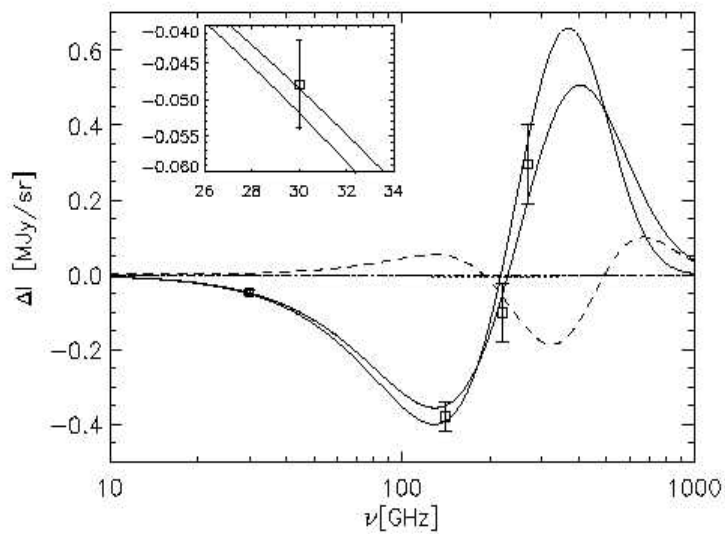


Data:  
LaRoque et al, astro-ph/0204134



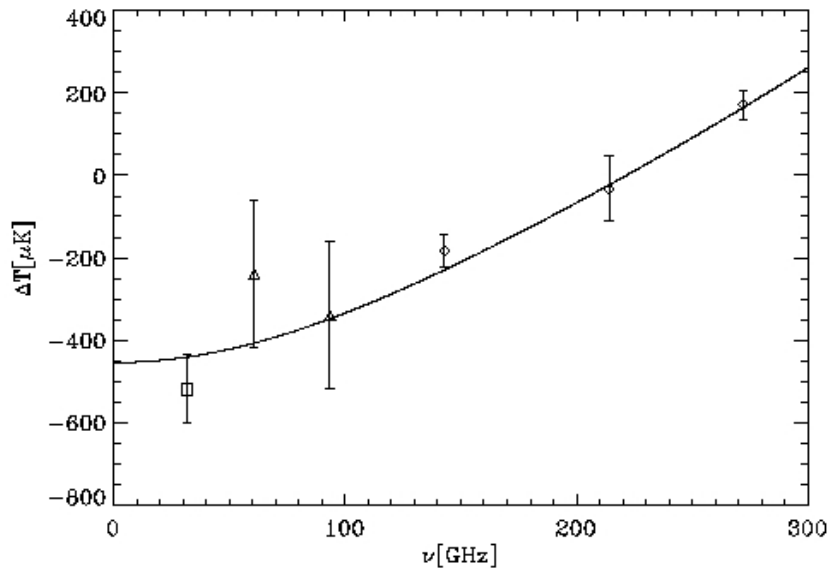
Temperature analysis:  
Hansen et al, ApJ 537 (2002) L69

# A2163 temperature through SZ, now in colour ☺



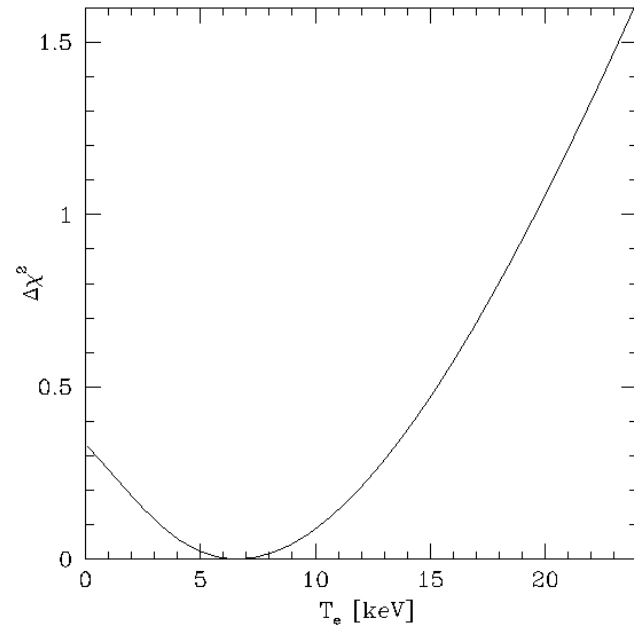
# Also Coma temperature determined

Result is  $T < 20$  keV at  $1\sigma$



Data:

Battistelli et al, ApJ, 598 (2003) L75



Temperature analysis:

Hansen, New Astron 9 (2004) 279

# What does the SZ formula look like

Non-relativistic:

$$\Delta I = \mathbf{y} * [\mathbf{g}(\mathbf{x})] - \beta * \tau * a(\mathbf{x})$$

Including relativistic effects:

$$\Delta I = \mathbf{y} * [\mathbf{g}(\mathbf{x}) + \mathbf{T} * \delta(\mathbf{x}, \mathbf{T})] - \beta * \tau * a(\mathbf{x})$$

Everything is integrated along the line of sight.

$\delta(\mathbf{x}, \mathbf{T})$  is almost independent of  $\mathbf{T}$  (Diego et al, MNRAS 338 (2003) 796)

The Compton parameter,  $y$ , is dominant

$$\Delta I = y * \{ [g(x) + T * \delta(x)] - \beta * \tau/y * a(x) \}$$

So if the cluster is **isothermal**, then this is simple, since

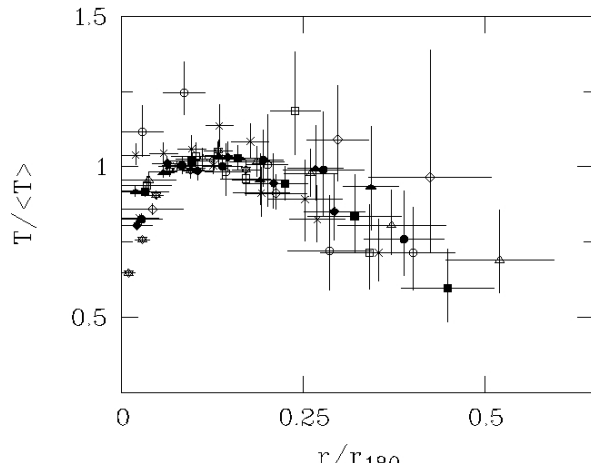
$$\tau/y \sim 1/T$$

$$\Delta I = y * \{ [g(x) + T * \delta(x)] - \beta / T * a(x) \}$$

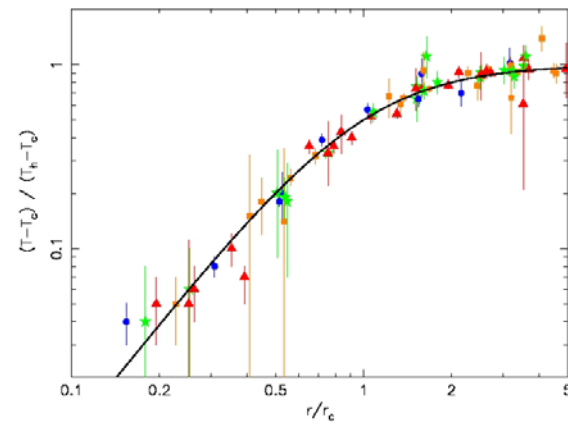
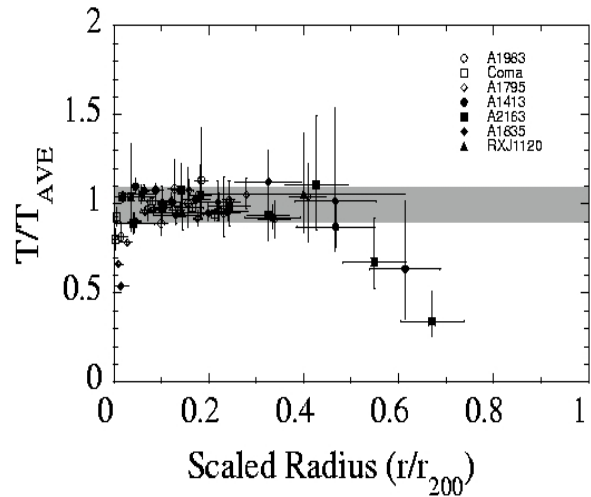
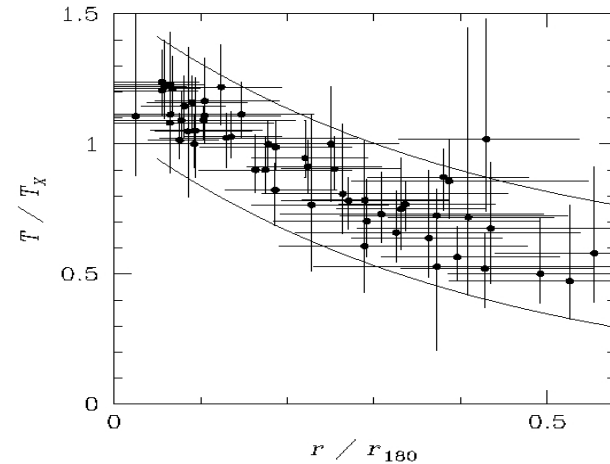
Where the **same**  $T$  appears twice

# However, the clusters are **not** isothermal

De Grandi & Molendi, ApJ 567 (2002) 1



Markevitch et al, ApJ 503 (1998) 77



Kaastra et al, A&A 413 (2003) 415

Arnaud et al, astro-ph/0312398

X-ray is measuring emission weighted quantities

Typically one uses emissivity  
 $\epsilon \sim n^2 * T$ , and one has

$$\langle O \rangle_X = \frac{\int O \epsilon dt}{\int \epsilon dt}$$

SZ is measuring Compton weighted quantities

$$\langle O \rangle_Y = \frac{\int O f_Y dt}{\int f_Y dt}$$

where  $f_Y$  is somehow related to the local  $y \sim n * T$

The formulae look simple after a few definitions

$$n_e(r) = n_e^0 f_{n_e}(r)$$

$$T_e(r) = T_e^0 f_{T_e}(r)$$

$$f_y = f_{n_e} f_{T_e}$$

$$y = \kappa_y \int f_y dl$$

$$\kappa_y = \frac{\sigma_T n_e^0 k T_e^0}{m_e c^2}$$

$$(T_e)_y = \frac{\int T_e f_y dl}{\int f_y dl} = \frac{\kappa_y \int T_e f_y dl}{y}$$

SZ temperature is **Compton weighted**

(Hansen, MNRAS 351 (2004) L5)

## A working example (1)

Consider a beta-model cluster,  $\beta = 0.61$

$$n_e(r) = n_e^0 \left( 1 + \left( \frac{r}{r_c} \right)^2 \right)^{-\beta 3/2}$$

and with a non-trivial temperature profile,  $\alpha = 0.56$

$$T_e(r) = T_e^0 \left( 1 + \frac{r}{r_c} \right)^{-\alpha}$$

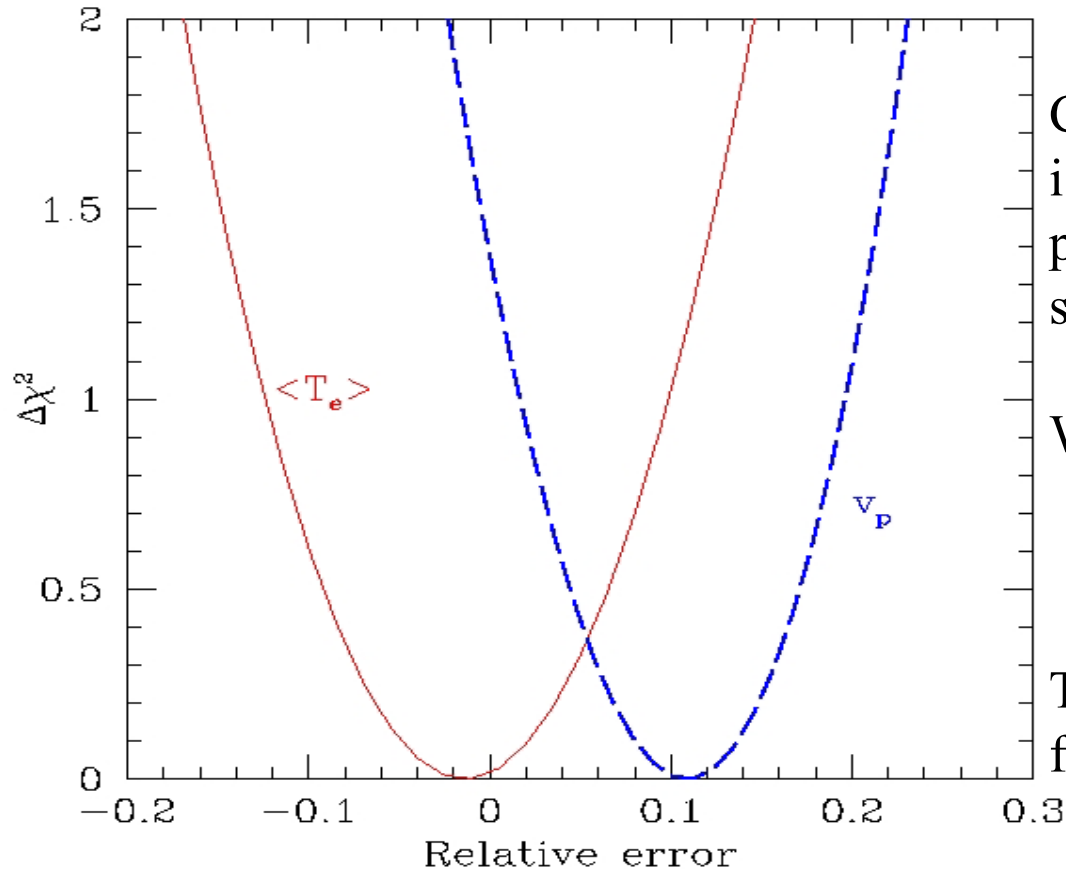
If central temperature is 10 keV, then the Compton averaged temperature is only 7.1 keV, and  $\langle 1/T_e \rangle \langle T_e \rangle = 1.12$

$$\Delta I = \mathbf{y} * \left\{ [g(\mathbf{x}) + \langle \mathbf{T} \rangle * \delta(\mathbf{x})] - \beta * \langle 1/\mathbf{T} \rangle * a(\mathbf{x}) \right\}$$

Now it is **not** the same  $\mathbf{T}$  that appears, since  $\langle \mathbf{T} \rangle$  and  $1/\langle 1/\mathbf{T} \rangle$  differ

## A working example (2)

We consider 4 observing frequencies at 90, 150, 220 and 270 GHz,  
and assume sensitivity of  $1\mu\text{K}$  for each channel



Compton weighted temperature  
is found correctly, however, the  
peculiar velocity is **systematically**  
shifted by 10-12%.

We had predicted it to be

$$\langle 1/T_e \rangle \langle T_e \rangle = 1.12$$

This is independent of observing  
frequencies and sensitivity.

## A working example (3)

The emission weighted temperature (from X-ray) for this cluster is about 8 keV. Using that would imply a systematic error on the peculiar velocity of about 20%.

Different (equally realistic) beta-models, or different radial temperature dependence lead to 10-20% systematic error on the peculiar velocity.

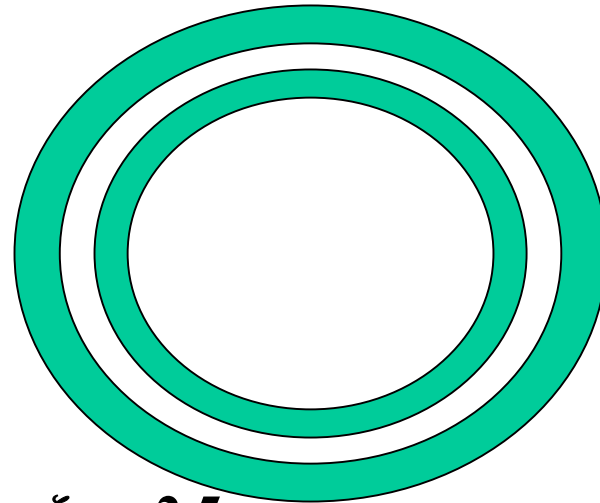
Peculiar velocities are most interesting for a large sample of clusters, for which we cannot expect to have accurate x-ray determinations of both density and temperature profiles.

Solution: **SZ onion peeling**

# SZ onion peeling

- First you observe and analyse outer ring
- Observe next ring, and subtract signal from outer rings
- Repeat

This works only if error-bars from outer rings don't induce too larger errors on inner rings



We can write  $n * T \sim r^{\zeta_y}$  Where  $\zeta_y \approx -2.5$

Compare to X-ray, where  $\varepsilon \sim r^{\zeta_x}$  and  $\zeta_x \approx -4$

Therefore, if we have resolution (and sufficient sensitivity) then the error-bars should remain under control

# Conclusions

- Clusters have a non-trivial temperature profile
- The SZ temperature is **Compton weighted**
  - Be extremely careful when using X-ray observed quantities in the SZ analysis
- The derived peculiar velocity is **systematically** shifted by 10-20%
- A nice solution is to achieve simultaneous good angular resolution and high sensitivity in order to deproject (onion peel) the SZ signal

**The End**