

The Metallicity Ladder in Resolved Stellar Populations



Manuela Zoccali (PUC / ESO) Ivo Saviane (ESO)





The (surface) chemical composition of a star

Theoretical scale:

Mass Fraction normalized to unity: X + Y + Z = 1



	~	•	
First stars in the Universe:	0.75 +	- 0.25	+ 0
Sun	0.71 +	+ 0.27	+ 0.02

Observational scale:

$$[Fe/H] = \log \frac{A(Fe)}{A(Fe)_{\odot}}$$
 with A(Fe) = N_{Fe} / N_H = Nr of Fe atoms / Nr of H atoms

Conversion:

 $[M/H] = \log \frac{Z}{Z_{\odot}} = [Fe/H]$ only if Fe traces all the metals





The (surface) chemical composition of a star

Theoretical scale:

Observational scale:

$$[Fe/H] = \log \frac{A(Fe)}{A(Fe)_{\odot}}$$

with A(Fe) = N_{Fe} / N_H = Nr of Fe atoms / Nr of H atoms

Conversion:

$$[M/H] = \log \frac{Z}{Z_{\odot}} = [Fe/H]$$
 only if Fe traces all the metals



Χ

Ζ

Υ



The spectrum of a star

Continuum (Black Body) + Fraunhofer Lines







What is a blackbody?

DEFINITION:

A container that is completely closed except for a very small hole in one wall. Any light entering the hole has a very small probability of finding its way out again, and eventually will be absorbed by the walls or the gas inside the container: this is a perfect absorber - all light that enters the hole is absorbed inside.



Eventually the photon finds the hole again and gets out, but this happens only after a lot of bouncing against the walls, i.e., after many interactions with the box material. This is the definition of "thermodynamic equilibrium", a condition that is also fulfilled in (most of) the stellar atmospheres.

By heating the box, one "heats" also the photons in it, and the resulting energy (=frequency) distribution of the outcoming photons also changes, according to the following laws:

Wien Law

 2.9×10^{-1}

Stefan-Boltzmann Law

$$F = \sigma T^4$$



Spectral Resolution

 $\frac{\lambda}{\Delta\lambda}$ R

Resolution is crucial in spectroscopy, not only to separate (resolve) individual spectral lines, but also to define continuum in crowded spectra (=metal rich stars, cold stars)





Spectral Resolution

R

Resolution is crucial in spectroscopy, not only to separate (resolve) individual spectral lines, but also to define continuum in crowded spectra (=metal rich stars, cold stars)





Spectral Resolution

R

Resolution is crucial in spectroscopy, not only to separate (resolve) individual spectral lines, but also to define continuum in crowded spectra (=metal rich stars, cold stars)





Chemical Abundance determination from HR spectra





Line profiles and curves of growth



Figure 9.20 Voigt profiles of the K line of Ca II. The shallowest line is produced by $N_a = 3.4 \times 10^{11}$ ions cm⁻², and the ions are ten times more abundant for each successively broader line. (Adapted from Novotny, *Introduction to Stellar Atmospheres and Interiors*, Oxford University Press, New York, 1973.)

For small optical depths $(\tau_{\upsilon} << I)$ it can be demonstrated that:

W/ λ = 8.85×10⁻¹³ N_i $f_{\lambda} \lambda$

where: Ni = column density of element f_{λ} = oscillator strength (sort of a

transition probability for that line)

[pag. 206 "Atomic Astrophysics and Spectroscopy" Pradhan & Nahar]



Figure 9.22 A general curve of growth for the Sun. (Figure from Aller, Atoms, Stars, and Nebulae, Revised Edition, Harvard University Press, Cambridge, MA, 1971.)



From EWs to Abundances

How many atoms of a given element can contribute to a given line?





From EWs to Abundances

How many atoms of a given element can contribute to a given line?

The problem is greatly simplified by the assumption of Local Thermodynamical Equilibrum:

1) The distribution of kinetic energies follows the **Maxwell law**

this tells you how many atoms change level due to collisions [they do not produce a line]

$$f(v) = 4\pi v^2 \left(\frac{m}{2\pi KT}\right)^{3/2} \exp\left(\frac{-mv^2}{2KT}\right)$$

2) The distribution of photon energy follows the **Planck's law**

$$B_{\nu}(T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT} - 1}$$

this tells you how photons of a given λ (or $\nu)$ are available

3) The distribution of electrons among different excitation states follows the **Boltzman equation** this tells you how many electrons there are in different atomic levels

$$\frac{N_i}{N_j} = \left(\frac{g_i}{g_j}\right) \exp\left(\frac{-(Ei - Ej)}{KT}\right)$$

4) The distribution of atoms among different ionization states follows the **Saha equation**

this tells you how many atoms are ionized

$$\log \frac{N^{I}}{N_{0}} = \log \frac{u^{I}}{u^{0}} + 2.5 \log T - \frac{5040}{T} \chi_{\text{ion}} - \log P_{e} - 0.176$$



The model atmosphere

gives T and P_e of the visible layers of the stellar atmosphere together with its **opacity**.





Bound-Bound absorption: small -except at those discrete wavelengths capable of producing a transition. i.e., responsible for forming absorption lines Bound-Free absorption: photoionisation -occurs when a photon has sufficient energy to ionize atom. The freed e⁻ can have any energy, thus this is a source of continuum opacity.

Free-Free absorption: a scattering process. A free electron absorbs a photon, causing the speed of the electron to increase. Can occur for a range of λ , so it is a source of continuum opacity.

Electron scattering: a photon is scattered, but not absorbed by a free electron. A very inefficient scattering process only really important at high temperatures -where it dominates



The iron distribution of a complex stellar population



In order to derive the Metallicity Distribution Function (MDF) of a SP it is necessary to select a target box that includes stars of all the metallicities present in the system.

One should also make sure that targets are \sim evenly distributed within the box.

+E\$+ 0 +

The iron distribution of a complex stellar population



The iron distribution of a complex stellar population





The Metallicity Distribution of the Galactic bulge



GaiaESO: Rojas-Arriagada et al. (2014, 2017)

ARGOS: Ness et al. (2013)



outer



The Metallicity Distribution of the Galactic bulge

